

Classical Mechanics

Phys105A, Winter 2007

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Formalities

- Latest news and course slides always found on the Phys105A site at [http://www.cs.ucsb.edu/~vandam/...](http://www.cs.ucsb.edu/~vandam/)
- Homework 3 has been posted.
It is due Monday January 29, 11:30 am.
Be sure to download the 'improved' version.
- Midterm is scheduled for Week 6, Thursday Feb 15.
The material will be Chapters 1–4; it is not open book, but you are allowed a 'cheat sheet'.

Force as a Gradient of U

Using the *gradient* we can summarize the connection between a potential U and its corresponding conservative force \mathbf{F} by the statement:

$$\mathbf{F} = -\nabla U$$

$$= -\frac{\partial U}{\partial x} \hat{\mathbf{e}}_x - \frac{\partial U}{\partial y} \hat{\mathbf{e}}_y - \frac{\partial U}{\partial z} \hat{\mathbf{e}}_z$$

Note that each possible potential U gives rise to a conservative force $\mathbf{F} = -\nabla U$.

The “Del” Operator ∇

You can think of del as $\nabla = \partial/\partial x + \partial/\partial y + \partial/\partial z$.

The *gradient* ∇U of a scalar field U is a vector field:

$$\nabla U = \partial U/\partial x \mathbf{e}_x + \partial U/\partial y \mathbf{e}_y + \partial U/\partial z \mathbf{e}_z.$$

The *curl* $\nabla \times \mathbf{V}$ of a vector field \mathbf{V} is a vector field:

$$\begin{aligned} \nabla \times \mathbf{V} = & (\partial V_z/\partial y - \partial V_y/\partial z) \mathbf{e}_x + (\partial V_x/\partial z - \partial V_z/\partial x) \mathbf{e}_y \\ & + (\partial V_y/\partial x - \partial V_x/\partial y) \mathbf{e}_z. \end{aligned}$$

The *divergence* $\nabla \cdot \mathbf{V}$ of a vector field \mathbf{V} is a scalar field:

$$\nabla \cdot \mathbf{V} = \partial V_x/\partial x + \partial V_y/\partial y + \partial V_z/\partial z.$$

$\nabla \times \mathbf{F} = \mathbf{0} \iff \mathbf{F}$ is Conservative

The fact that \mathbf{F} is a conservative force is equivalent with the statement that $\nabla \times \mathbf{F} = \mathbf{0}$ at all positions (x, y, z) .

To prove this we use Stokes' Theorem:

$$\oint_{\Gamma} \mathbf{F} \cdot d\mathbf{r} = \int_{\Sigma} \nabla \times \mathbf{F} \cdot d\boldsymbol{\sigma}$$

where on the left hand we integrate over the boundary Γ of the surface Σ , while on the right hand we integrate over the surface itself (using the normal vector $d\boldsymbol{\sigma} = \mathbf{e}_n \cdot d\sigma$).

Understanding Stokes' Thm

To understand Stokes' theorem it helps to know that for small squares $(\mathbf{r}, \mathbf{r}+\Delta\mathbf{x}, \mathbf{r}+\Delta\mathbf{y}, \mathbf{r}+\Delta\mathbf{x}+\Delta\mathbf{y})$ we have at \mathbf{r} :

$$\nabla \times \mathbf{F} = \frac{1}{\Delta x \Delta y} \oint_{\mathbf{r} \rightarrow \mathbf{r}+\Delta\mathbf{x} \rightarrow \mathbf{r}+\Delta\mathbf{x}+\Delta\mathbf{y} \rightarrow \mathbf{r}+\Delta\mathbf{y} \rightarrow \mathbf{r}} \mathbf{F} \cdot d\mathbf{r}' \quad \text{as } \Delta x, \Delta y \rightarrow 0.$$

By patching such infinitesimal squares we can make proper surfaces, thus obtaining Stokes' Theorem.

The requirement $\nabla \times \mathbf{F} = \mathbf{0}$ certifies that each (infinitesimal) circular path will have a net work $W(\mathbf{r} \rightarrow \mathbf{r}) = 0$.

Time Dependent U

- For time dependent U, we have at any given moment that $\mathbf{F} = -\nabla U$ is conservative, yet the total energy $E = T+U$ is not conserved.
- $dT = m(\mathbf{v} \cdot d\mathbf{v}) = m(\mathbf{v} \cdot \mathbf{a})dt = \mathbf{F} \cdot d\mathbf{r}$
- $dU = \frac{\partial U}{\partial x} dx + \frac{\partial U}{\partial y} dy + \frac{\partial U}{\partial z} dz + \frac{\partial U}{\partial t} dt$
 $= \nabla U \cdot d\mathbf{r} + \frac{\partial U}{\partial t} dt = -\mathbf{F} \cdot d\mathbf{r} + \frac{\partial U}{\partial t} dt$
- So, $dE = d(T+U) = \frac{\partial U}{\partial t} dt$ and $\Delta E = \int \frac{\partial U}{\partial t} dt$.