# *Differential Equations Basics*



## *Differential Equations*

- **Equations with derivative symbols :-)**
- Ordinary differential equations
	- $\Box$  initial value problems
- Partial differential equations
	- **□** boundary value problems, initial-boundary value problems
- **Exact solutions** 
	- **Q** guess work
- Numerical solutions
	- **Q** guess work (but computer does it!)
- $\frac{1}{2}$ *Warning: This is not a course on numerical analysis! Take an appropriate course if interested*



*Ordinary Differential Equations*  Usually the independent variable is treated as time  $(t)$ (*t*)  $\dot{x} = f(x,t)$ Example





## *More Examples*

$$
\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = -k \begin{bmatrix} \cos(xt) \\ \sin(yt) \end{bmatrix}
$$





## *Intuition*

- **Imagine a vector field** 
	- **T** field (flow) direction specified by
		- location, and

 $\triangleright$  time

- **a** a particle in the vector field will drift according to the current flow direction where the particle is
- $\Box$  time independent dropping something in a river and watching it flows
- $\Box$  time dependent dropping something in the ocean and watching ocean waves carry it



## *Numerical Solutions*

- Advance a particle from its initial location through (discrete) time steps
- $\cdot$  Many issues
	- **accuracy**
	- **<u>a</u>** stability
	- $\Box$  efficiency
- Man y solutions
	- **□** Euler (explicit, implicit, modified)
	- Midpoint and Runge-Kutta
	- **□** Multistep, predictor-corrector
	- **Adaptive step size**





### *Issues*

#### **Accuracy**

 $\Box$  How close is the discrete solution going to be to the continuous solution?

**❖ Stability** 

**Numerical error unavoidable** 

 $\Box$  Error accumulation to grow or decay?

- **\*** Efficiency
	- $\Box$  How fast for each time step?
	- **□** How large can each time step be?



## *(Explicit) Euler Method*

• Go in the tangent direction with distance controlled by time step





## *Questions*

**Euler method accurate?**  $\cdot$  Is Euler method stable? **★ How close are the the exact and Euler** solutions?

Exact solution

Euler solution

 $\dot{x} = -kx$ 

*x* ٥  $\dot{x} = -kx$   $\dot{x} =$  $kn\Delta t$  $\Rightarrow$  *x*(*n* $\Delta t$ ) = *x*<sub>o</sub>  $e^{-kn\Delta t}$ *kt*  $\Rightarrow x = x_o e^-$ 

 $(n\Delta t) = x_0 e^{-kn\Delta t}$   $= ... = (1 - k\Delta t)^n x_0$  $x_n = x_{n-1} - k\Delta t x_{n-1} = (1 - k\Delta t)x_{n-1}$  $-k\Delta t x_{n-}$  $=(1-k\Delta t)x_{n-}$ 



#### $(\dot{x} = -x)$   $x_n = (1 - \Delta t)^n x_0$  30 - second simulation





## *A graphical interpretation*

 $\dot{x} = -kx$ At large step sizes At even large step sizes





### *Observations (Explicit Euler)*

- Explicit Euler method can be made accurate as long as you can tolerate (*very*) small step size
- It can be made stable (not blowing up!) with (*appropriate*) small step size
- **Explicit Euler is easy to understand**
- **★ But it is seldom used in the real world (homework** assignments are *not* real world, but should be treated as such  $\circledcirc$  )  $\dot{x} = -kx$

$$
x_{n+1} = (1 - k\Delta t)^n x_o
$$
  
\n
$$
\Rightarrow 1 - k\Delta t \ge -1
$$
  
\n
$$
\Rightarrow k\Delta t \le 2
$$





## *How Small Should Step Size Be?*

• The larger the k, the smaller the step size  $\cdot$  Equation of large k, or system of equations with large varying k, are stiff

$$
\begin{aligned}\n\dot{x} &= -kx \\
x_{n+1} &= (1 - k\Delta t)^n x_o \\
\Rightarrow 1 - k\Delta t &\geq -1 \\
\Rightarrow k\Delta t &\leq 2 \\
\Rightarrow \Delta t &\leq \frac{2}{k}\n\end{aligned}
$$



#### *(Implicit) Euler Method*

 $\dot{x} = f(x,t)$  $= f(x,t)$  $\frac{dx}{dt} = f(x,t)$   $\frac{dx}{dt} = -kx$  $(x_{n+1}, t_{n+1})$  $\cdot$  $+1$ <sup>2</sup> $n+$  $\frac{1}{\Delta t} =$ *n n n n*  $f(x_{n+1}, t)$ *t*  $x_{n+1} - x$ *x* $\lambda_{n+1} = x_n + \Delta t \cdot f(x_{n+1}, t_{n+1})$  $_{+1}-\Delta t\cdot f(x_{n+1},t_{n+1})$  $x_{n+1} = x_n + \Delta t \cdot f(x_{n+1}, t_n)$  $\Lambda_{n+1} - \Delta t \cdot f(x_{n+1}, t_{n+1}) = x_n$  $x_{n+1} - \Delta t \cdot f(x_{n+1}, t_{n+1}) = x$ *nnxx*

$$
\begin{aligned}\n\dot{x} &= -kx \\
\frac{dx}{dt} &= -kx \\
\frac{x_{n+1} - x_n}{\Delta t} &= -kx_{n+1} \\
x_{n+1} + k\Delta tx_{n+1} &= x_n \\
x_{n+1} &= \frac{1}{1 + k\Delta t} x_n\n\end{aligned}
$$



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## *Observations (Implicit Euler)*

- Accuracy depends on step size
- Absolute stability, regardless of the step size
- A magic bullet for many "stiff" system
- More complicated (not necessarily direct evaluation), try  $f = ax^2 + bx + c$  or  $f = ax^2 + bx +$

 $f = a \cos x + b \sin x$ 

**Exercication maybe needed Expensive, seldom used in the real world** (think about system of equations)



## *(Modified) Euler Method*





## *Observations (Modified Euler)*

- Use trapezoidal rule
- Accuracy depends on step size (Surprisingly: it is more accurate than either explicit or implicit method)
- *Absolute* stability, *regardless of* the step size
- More complicated (not necessarily direct evaluation)
- **★ Basis for other advanced techniques**



## *Accuracy of Euler Methods*

$$
x(t+\Delta t) = x(t) + \dot{x}(t)\Delta t + \ddot{x}(t)\frac{\Delta t^2}{2!} + \ddot{x}(t)\frac{\Delta t^3}{3!} + \cdots
$$

Explicit Euler:  $x(t + \Delta t) = x(t) + \dot{x}(t)\Delta t$  $\frac{d}{dt}$  *Error*:  $O(\Delta t^2)$ 

 $I$ mplicit Euler :  $x(t + \Delta t) = x(t) + \dot{x}(t + \Delta t) \Delta t$ 

 $(t + \Delta t) = \dot{x}(t) + \ddot{x}(t)\Delta t + \dddot{x}(t) - \frac{1}{2!}$  $\dot{x}(t + \Delta t) = \dot{x}(t) + \ddot{x}(t)\Delta t + \dddot{x}(t)\frac{\Delta t^2}{2} +$  $+\Delta t$ ) =  $\dot{x}(t) + \ddot{x}(t)\Delta t + \dddot{x}(t)\frac{\Delta t}{t}$  $\dot{x}(t+\Delta t) = \dot{x}(t) + \ddot{x}(t)\Delta t + \dddot{x}(t) - + \cdots$ 

$$
x(t+\Delta t) = x(t) + \dot{x}(t)\Delta t + \ddot{x}(t)\Delta t^{2} + \dddot{x}(t)\frac{\Delta t^{3}}{2!} + \cdots
$$

 $Error$  :  $O({\Delta t}^2)$ 



*d*

*dt*

$$
Accuracy of Euler Methods (cont.)
$$
  
Modified Euler:  $x(t+\Delta t) = x(t) + \frac{1}{2}(\frac{x(t) + \dot{x}(t+\Delta t))\Delta t}{2!}$   

$$
\dot{x}(t+\Delta t) = \dot{x}(t) + \ddot{x}(t)\Delta t + \ddot{x}(t)\frac{\Delta t^2}{2!} + \cdots
$$
  

$$
x(t+\Delta t) = x(t) + \frac{1}{2}\dot{x}(t)\Delta t + \frac{1}{2}\{\dot{x}(t)\Delta t + \ddot{x}(t)\Delta t^2 + \ddot{x}(t)\frac{\Delta t^3}{2!} + \cdots\}
$$
  

$$
= x(t) + \dot{x}(t)\Delta t + \frac{1}{2}\ddot{x}(t)\Delta t^2 + \frac{1}{2}\ddot{x}(t)\frac{\Delta t^3}{2!} + \cdots
$$
  
*Error*:  $O(\Delta t^3)$ 



#### *Why Implicit Methods Don't Blow Up?*

 $\dot{x} = -kx, k > 0$ 



## *Stability=(?)Accuracy*

- In general, you can get a stable solution by using <sup>a</sup> reasonably small step size
- To get high accuracy, step size needs to be very small (with dumb methods)
- There are other techniques to increase accurac y at lar ge step size



## *Why not Higher Orders?*

#### $\cdot$  If

- **T** the goal is to reduce the *local* approximation error
- **Local approximation error is expressed in terms** of Taylor series
- **\*** Then
	- Why not compute more terms in the expansion?  $\ddot{x}(t)\Delta t + \ddot{x}(t)$   $\frac{d}{dx}$  +  $\dddot{x}(t)$   $\frac{d}{dx}$  +  $+\ddot{x}(t)\frac{\Delta}{ }$  $(t + \Delta t) = x(t) + \dot{x}(t)\Delta t + \ddot{x}(t)\frac{\Delta t^2}{\Delta t} + \dddot{x}(t)$  $\frac{2}{1+x}(t) \frac{\Delta t^3}{t^2}$  $x(t + \Delta t) = x(t) + \dot{x}(t)\Delta t + \ddot{x}(t)\frac{\Delta t^2}{2!} + \dddot{x}(t)\frac{\Delta t^3}{3!} + \cdots$  $+\cdots$ +  $f''\frac{\Delta}{ }$  $(t + \Delta t) = x(t) + f\Delta t + f'\frac{\Delta t^2}{2!} + f'''\frac{\Delta t}{3!}$  $\frac{2}{1} + f'' \frac{\Delta t^3}{2}$  $x(t + \Delta t) = x(t) + f\Delta t + f' \frac{\Delta t}{2}$  $\dot{x} = f(x,t)$



### *Answer*

- $\cdot$  Only  $f(x,t)$  (or x') is given explicitly  $\mathbf{f}'(x,t)$ , etc (*or x*<sup>"</sup>, x<sup>"</sup>) can be hard to evaluate
- There are other methods to match the terms in Taylor series to achieve high-order of accuracy
	- **□ Single-step, Compounding methods**
	- **□** Mutli-step methods



*Midpoint Method (2nd order Runge Runge-Kutta)* • Compound one-step method  $\cdot$  *Explicit* for easy computing **★ 2nd order is very similar to trapezoidal rule** • Caution: there are a number of 2<sup>nd</sup> order Runge-Kutta formulation, all with the same error bound



## *Trapezoidal*

Notice

\nModel Euler

\n
$$
\begin{aligned}\n\text{where } \hat{x} &= f(x, t) \\
\frac{dx}{dt} &= f(x, t) \\
\frac{x_{n+1} - x_n}{\Delta t} &= \frac{1}{2} (f(x_n, t_n) + f\left(\frac{x_n + \Delta t f(x_n, t_n)}{x_n + \Delta t f(x_n, t_n)}\right), t_{n+1}) \\
\frac{dx}{dt} &= f(x, t) \\
\frac{x_{n+1} - x_n}{\Delta t} &= \frac{1}{2} (f(x_n, t_n) + f\left(\frac{x_{n+1}}{x_{n+1}}\right), t_{n+1}) \\
\frac{x_{n+1} - x_n}{\Delta t} &= \frac{1}{2} (f(x_n, t_n) + f\left(\frac{x_{n+1}}{x_{n+1}}\right), t_{n+1}) \\
\frac{E_x}{dt}\n\end{aligned}
$$



*x*

*Example*





## *Facts*

- $\bullet$   $O(\Delta t^3)$  accuracy
	- $\Box$  Explicit, easy to compute, but not absolute stable
- Other compounding methods possible • Compound can be done many many times,
- the most famous is the 4th order Runge-Kutta with  $O(\Delta t^5)$



## *General Compounding Methods*

- $\triangleleft$  Evaluate the expression  $f(x,t)$  *multiple* times in <sup>a</sup> *single* step
	- At *different* location ( *<sup>x</sup>*) and *different* time ( *t* )
	- $\Box$  But all *explicit*
- **→ 2<sup>nd</sup> order (midpoint) actually corresponds to** determine a, b,c in the following expression  $\Box$  Evaluate the derivative at different place (b), different time  $(c)$ , and different weighting  $(a)$

 $\dot{x} = af(x+b, t+c)$ )  $cf \quad \dot{x} = f(x,t)$  for Explict Euler



$$
\mathbf{Matching Process}
$$
\n
$$
x_{n+1} = x_n + \Delta t x_n + \frac{\Delta t^2}{2!} x_n^* + \cdots
$$
\n
$$
x_{n+1} = x_n + \Delta t f(x_n, t_n) + \frac{\Delta t^2}{2!} f'(x_n, t_n) + \cdots
$$
\n
$$
= x_n + \Delta t f(x_n, t_n) + \frac{\Delta t^2}{2!} \left( \frac{\partial f(x_n, t_n)}{\partial t} + \frac{\partial f(x_n, t_n)}{\partial x} \right) + \cdots
$$
\n
$$
= x_n + \Delta t f(x_n, t_n) + \frac{\Delta t^2}{2!} \left( \frac{\partial f(x_n, t_n)}{\partial t} + \frac{\partial f(x_n, t_n)}{\partial x} \right) f(x_n, t_n) + \cdots
$$
\n
$$
= x_n + \Delta t f(x_n, t_n) + \frac{\Delta t^2}{2!} \frac{\partial f(x_n, t_n)}{\partial t} + \frac{\Delta t^2}{2!} \frac{\partial f(x_n, t_n)}{\partial x} f(x_n, t_n) + \cdots
$$
\n
$$
= x_n + \Delta t f(x_n, t_n) + \frac{\Delta t^2}{2!} \frac{\partial f(x_n, t_n)}{\partial t} + \frac{\Delta t^2}{2!} \frac{\partial f(x_n, t_n)}{\partial x} f(x_n, t_n)
$$

$$
x_{n+1} = x_n + a\Delta tf(x+b, t+c)
$$
  
\n
$$
x_{n+1} = x_n + a\Delta t[f(x, t) + \frac{\partial f}{\partial x}b + \frac{\partial f}{\partial t}c + \cdots]
$$
  
\n
$$
x_{n+1} = x_n + \Delta taf(x, t) + \Delta t \frac{\partial f}{\partial x}ab + \Delta t \frac{\partial f}{\partial t}ac + \cdots]
$$



*Matching Process*



#### *General Compound Methods*

 Successive recursions (many more times)  $\cdot$  4<sup>th</sup>-order is the most popular



$$
f_4 = f(x + \Delta t f_3, t + \Delta t)
$$
  
\n
$$
f_3 = \Delta t f(x + \frac{\Delta t}{2} f_2, t + \frac{\Delta t}{2})
$$
  
\n
$$
f_2 = \Delta t f(x + \frac{\Delta t}{2} f_1, t + \frac{\Delta t}{2})
$$
  
\n
$$
f_1 = \Delta t f(x, t)
$$
  
\n
$$
x_{n+1} = x_n + \frac{\Delta t}{6} (f_1 + 2f_2 + 2f_3 + f_4)
$$
  
\nOne evaluation on the left  $\uparrow$  Two evaluations in the middle  
\nOne evaluation on the right



## *Multi-Step Methods*

- Utilize multiple *pas<sup>t</sup>* and *presen<sup>t</sup>* values instead of just one (at time *n - 1* )
- **Explicit, easy to compute**
- Implicit also possible for stability
- Can be made into arbitrary orders of accurac y (usin g enou gh terms)
- $\cdot$  For smooth fields  $(f', f'')$ , etc. should exist)  **Compare to**



## *Comparison*

- **❖ Compound methods**
- *Multiple* evaluations of *f* in a single step
- $\bullet$  Do not use past values  $\bullet$
- Do not assume smoothness over time
- **❖ Good accuracy (4th** order)

**Multi-step methods**  *Single* evaluation of *f* in a single step Use pas<sup>t</sup> values of *f*

**Work best with** ↓ s over time smooth fields  $(f', f'',$ <br>racy  $(4^{\text{th}} \t f'')$ , etc. exist)

 $\cdot$  Very high accuracy



## *Multi-Step Methods*

 Explicit (Adams-Bashforth)

#### Implicit (Adams-Moulton)

- *f* at previous times Already calculated, might as well use them
- $1 J n$   $2 J n-1$   $\cdots$   $v_{p} J n-p+1$  $\frac{1}{\Lambda} \frac{X_n}{\Lambda} = c_1 f_n + c_2 f'_{n-1} + \ldots + c_p f_{n-p+1}$  $\frac{1}{\Delta t} = c_1 f_n + c_2 f_{n-1} + ... +$ F. *n p p n*-*p*  $c_1 f_n + c_2 f_{n-1} + ... + c_n f_n$  $x_{n+1} - x$ *t*  $\mathbf{r}$  *r*  $\mathbf{r}$  *<i>r*  $\mathbf{r}$  *r*  $\mathbf{r}$  *r*  $\mathbf{r}$  *r*

$$
\dot{x} = f(x, t)
$$
\n
$$
\frac{x_{n+1} - x_n}{\Delta t} = \underbrace{c_0 f_{n+1}} + c_1 f_n + c_2 f_{n-1} + \dots + c_p f_{n-p+1}
$$

• Q: What should the coefficient c\_p be? • A: Chosen to match as many terms in the Taylor expansion a s poss i b l e



 $\dot{x} = f(x,t)$ 

#### *Example - Explicit, 2nd order*

$$
\underbrace{f_{n-1} = f_n - f_n' \Delta t}_{\Delta t} + \underbrace{f_n'' \Delta t}_{2!} = \underbrace{f_{n-1} - f_n' \Delta t}_{\Delta t} + \underbrace{f_n'' \Delta t}_{2!}
$$

 $\gamma$ 

 $\lambda$ 

$$
x_{n+1} = x_n + \alpha \Delta t f_n + \beta \Delta t \{f_n - f'_n \Delta t + f_n'' \frac{\Delta t^2}{2!} - \cdots \}
$$

$$
x_{n+1} = x_n + (\alpha + \beta) f_n \Delta t - \beta f_n' \Delta t^2 + \beta f_n'' \frac{\Delta t^3}{2!} - \cdots
$$

 $\alpha + \beta = 1$  $Error: O({\Delta t}^3)$  $\overline{2}$ 1 $\beta$ **TELE** 

 $\, + \,$ 

 $x_{n+1} - x_n$ 

 $\Delta t$ 

 $\frac{\overline{\lambda}_n}{\sqrt{n}} = \alpha f$ 



*Computer Graphics*

. . .

## *Observation:Multi-Steps*

#### **Again**

- $\Box$  implicit methods afford high accuracy with good stability
- $\Box$  the cost is again in evaluating the extra term in the implicit methods (pretty difficult if *f* is not linear!)
- **O** Good for smooth functions requiring high accuracy



#### **Explicit**



$$
\Delta t \leq \frac{stability\,threhold}{-k}
$$



#### *Predictor-Corrector Method*

 $\cdot$  Use the explicit formula to *predict* a new  $x_{n+1}$ <sup>\*</sup>  $x_{n+1}$  $x_{n+1}^* = x_n + \Delta t (c_1 f_n + c_2 f_{n-1} + ... + c_p f_{n-p+1})$  $\cdot$  Use  $_{x_{n+1}}$  to evaluate the right side  $(t_{n+1}, x_{n+1}^{*})$ \*  $f_{n+1} = f(t_{n+1}, x_{n+1})$ 

Use  $f_{n+1}$ <sup>\*</sup> in the implicit formula to correct the new  $x_{n+1}$ 

 $1 \cdot \nu_{n+1}$ 

$$
x_{n+1} = x_n + \Delta t (c_0 f_{n+1}^* + c_1 f_n + c_2 f_{n-1} + \dots + c_p f_{n-p+1})
$$



### *Error Sources*

- **So far we concentrated on truncation error** i.e., we didn't use enough terms in Taylor series
- $\cdot$  Another source is numerical (10^-16 for double due to limited machine precision)

$$
E = E_{round} + E_{truncated}
$$

 Eround: error in calculating derivatives Etruncate: omission in Taylor expansion



## *Error Sources*

- **Reduce step size** 
	- **O** Truncation error becomes smaller
	- **Numerical error** accumulates faster from many steps
	- **Simulation becomes** smaller
- **Increase step size** 
	- **T** Truncation error becomes larger
	- **Numerical error** accumulates slower from fewer steps
	- **Simulation becomes** faster

 $E$   $=$   $E_{round}$   $+$   $E_{truncated}$ 



## *Error Sources (cont)*

• Obviously, some trade-off exists with an optimial step size  $\sim$ (Eround/max(Etruncate)) $\sim$ 0.5  $\cdot$  Easier said than done **How to find Etruncate Example 12** Etruncate is not a constant over the whole domain Adaptive step size is the key



## *Adaptive Step Size*

• To speed up simulation needs large step size • To maintain stability and improve accuracy need small step size

• The step size should be adapted depending on the required accuracy





- Error can be estimated from the difference of two Euler steps
- Steps can be enlarged or shrunk based on the desired accuracy



## *Adaptive step sizes*





## *Last Word*

- If you are a numerical analyst, this does not apply to you  $\odot$
- *❖* Otherwise
	- **Don't use Explicit Euler (not accurate, blowing** up)
	- **D** Don't use Implicit Euler (too expensive)
	- $\square$  2<sup>nd</sup> or 4<sup>th</sup> order Runge-Kutta with step control is for you



## *Solver Structure*

- Heart is dxdt (a function you supply)
- Wrapped around by fixed step Euler, Midpoint, Runge-Kutta, etc.
- Wrapped around b y a step size controller





## *Example: 4th-order Runge-Kuta*

#include "nrutil.h"

```
void rk4(float y[], float dydx[], int n, float x, float h, float yout[],
    void (*derivs)(float, float [], float []))
Given values for the variables y [1...n] and their derivatives dydx [1...n] known at x. use the
fourth-order Runge-Kutta method to advance the solution over an interval h and return the
incremented variables as yout [1..n], which need not be a distinct array from y. The user
supplies the routine derivs (x, y, dydx), which returns derivatives dydx at x.
\mathcal{F}int i:
    float xh, hh, h6, *dym, *dyt, *yt;
    dym=vector(1, n);dyt = vector(1, n);
     vt=vector(1, n);
     hh=h*0.5;
     h6=h/6.0;
     xh = x + h h;
     for (i=1; i<=n; i++) yt[i]=y[i]+hh*dydx[i];First step.
     (*derivs)(xh,yt,dyt);
                                                           Second step.
     for (i=1; i<=n; i++) yt[i]=y[i]+hh*dyt[i];(*derivs)(xh,yt,dyn):Third step.
     for (i=1:i<=n:i++) {
         yt[i]=y[i]+h*dym[i];dyn[i] += dyt[i];Ł
     (*derivs)(x+h,yt,dyt);Fourth step.
                                                           Accumulate increments with proper
     for (i=1:i<=n:i++)yout [i] = y[i] + h6*(dydx[i] + dyt[i]+2.0*dym[i]);weights.
     free\_vector(yt,1,n):
     free\_vector(dyt, 1, n):
     free\_vector(dym, 1, n):
```
 $y_{n}$  $y_{n+1}$ h*Computer Graphics*

## *Example: RK with Step Control*

#include "nrutil.h" #define SAFETY 0.9 #define PCROW -0.2 #define PSHRNK -0.25 #define ERRCON 1.89e-4 The value ERRCON equals (5/SAFETY) raised to the power (1/PGROW), see use below. void rkqs(float y[], float dydx[], int n, float \*x, float htry, float eps, float yscal[], float \*hdid, float \*hnext, void (\*derivs) (float, float  $[]$ , float  $[]$ )) Fifth-order Runge-Kutta step with monitoring of local truncation error to ensure accuracy and adjust stepsize. Input are the dependent variable vector  $y[1, n]$  and its derivative dydx  $[1, n]$ at the starting value of the independent variable x. Also input are the stepsize to be attempted htry, the required accuracy eps, and the vector yscal[1..n] against which the error is scaled. On output, y and x are replaced by their new values, hdid is the stepsize that was actually accomplished, and hnext is the estimated next stepsize. derivs is the user-supplied routine that computes the right-hand side derivatives. -f void rkck(float y[], float dydx[], int n, float x, float h, float yout[], float yerr[], void (\*derivs)(float, float [], float [])); int i: float errmax, h, htemp, xnew, \*yerr, \*ytemp;  $\text{verr=vector}(1,n)$ ;  $y$ temp=vector $(1,n)$ ; Set stepsize to the initial trial value.  $h=htry;$ for  $(::)$  {  $rkck(y, dydx, n, *x, h, vtemp, verr, derivs)$ : Take a step.  $errmax=0.0$ : Evaluate accuracy. for  $(i=1; i<=n; i++)$  errmax=FMAX(errmax,fabs(yerr[i]/yscal[i])); Scale relative to required tolerance.  $errmax$  /=  $eps$ : Step succeeded. Compute size of next step. if (errmax  $\leq 1.0$ ) break; htemp=SAFETY\*h\*pow(errmax,PSHRNK); Truncation error too large, reduce stepsize.  $h = (h > = 0.0 ? FMAX(htemp, 0.1*h)$  :  $FMIN(htemp, 0.1*h))$ ; No more than a factor of 10.  $xnew = (*x) + h;$ if (xnew ==  $*x$ ) nrerror("stepsize underflow in rkqs"); ŀ. if (errmax > ERRCON) \*hnext=SAFETY\*h\*pow(errmax.PGROW): No more than a factor of 5 increase. else \*hnext=5.0\*h;  $*x \leftarrow (*hdid=h)$ : for  $(i=1; i<=n; i++)$   $y[i]=ytemp[i];$  $free\_vector(ytemp, 1, n)$ ; free\_vector(yerr,1,n);

y n $y_{n+1}$  $\rm h_{try}$  $y_{n+1}$ y n $\rm h_{did}$ *Computer Graphics*  $h_{\text{next}}$ 

#### *Example: Other Options*

void rkck(float  $y[]$ , float dydx $[]$ , int n, float x, float h, float yout $[]$ , float yerr[], void (\*derivs)(float, float  $[]$ , float  $[]$ )) Given values for n variables  $y[1..n]$  and their derivatives  $dydx[1..n]$  known at x, use the fifth-order Cash-Karp Runge-Kutta method to advance the solution over an interval h

and return the incremented variables as yout  $[1\, \ldots n]$ . Also return an estimate of the local truncation error in yout using the embedded fourth-order method. The user supplies the routine  $derivs(x, y, dydx)$ , which returns derivatives dydx at x.

![](_page_52_Figure_3.jpeg)

 Will advance by *h* (may take multiple steps, or multiple calls to rkqs)

![](_page_52_Picture_5.jpeg)

#### *Example: Other Options*

#include <math.h> #include "nrutil.h" #define MAXSTP 10000 #define TINY 1.0e-30

extern int kmax, kount;

extern float \*xp,\*\*yp,dxsav;

User storage for intermediate results. Preset kmax and dxsav in the calling program. If kmax  $\neq$  $0$  results are stored at approximate intervals dxs av in the arrays  $xp[1, .kount], yp[1, .nvar]$ [1..kount], where kount is output by odeint. Defining declarations for these variables, with memory allocations  $xp[1..kmax]$  and  $yp[1..nwar][1..kmax]$  for the arrays, should be in the calling program.

void odeint(float ystart[], int nvar, float x1, float x2, float eps, float h1, float hmin, int \*nok, int \*nbad, void (\*derivs)(float, float [], float []), void (\*rkqs)(float [], float [], int, float \*, float, float, float [], float \*, float \*, void (\*) (float, float  $[]$ , float  $[]$ )))

Runge-Kutta driver with adaptive stepsize control. Integrate starting values ystart [1..nvar] from x1 to x2 with accuracy eps, storing intermediate results in global variables. h1 should be set as a guessed first stepsize, hmin as the minimum allowed stepsize (can be zero). On output nok and nbad are the number of good and bad (but retried and fixed) steps taken, and ystart is replaced by values at the end of the integration interval. derivs is the user-supplied routine for calculating the right-hand side derivative, while rkqs is the name of the stepper routine to be used.

 $\rm h_2$ 

 $\mathbf{h}^{\cdot}_1$ 

 $\rm y_{end}$ 

hn

![](_page_53_Picture_8.jpeg)

## *Stiff Systems*

- Nightmare in numerical solutions
- Need small steps to maintain numerical stability
- But the interesting phenomena are really at large time scale
- \* Take a long time vs. blowing up
- May need implicit methods to achieve fast simulation without blowing up

![](_page_54_Picture_6.jpeg)

*Example*

$$
\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}
$$
  

$$
\begin{bmatrix} x_n \\ y_n \end{bmatrix} = \begin{bmatrix} x_{n-1} \\ y_{n-1} \end{bmatrix} + \Delta t \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x_{n-1} \\ y_{n-1} \end{bmatrix} = (I + \Delta t A) \begin{bmatrix} x_{n-1} \\ y_{n-1} \end{bmatrix}
$$
  
= ... =  $(I + \Delta t A)^n \begin{bmatrix} x_0 \\ y_0 \end{bmatrix}$ 

If A has eigen vect ors  $(u_1, u_2)$ , and eigen valu es  $(\lambda_1, \lambda_2)$  $=c_1u_1 + c_2u_2$  $x_{\rm o}$  $\Rightarrow$  *I* +  $\Delta t$ *A* has eigen vect ors ( $u_1, u_2$ ), and eigen valu es ( $1 + \Delta t \lambda_1$ ,  $1 + \Delta t \lambda_2$ )  $\Rightarrow \left[\begin{array}{c} x_0 \\ x_0 \end{array}\right] = c_1 u_1 +$  $(I + \Delta t A)^n (c_1 u_1 + c_2 u_2) = c_1 (1 + \Delta t \lambda_1)^n u_1 + c_2 (1 + \Delta t \lambda_2)^n u_2$  $\begin{bmatrix} 0 \\ y_0 \end{bmatrix}$ *y*  $x_n$   $(x_{n+1}, x_n)$ *n*  $|n| = (I + \Delta)$  $\Rightarrow$   $\begin{bmatrix} x_n \\ y_n \end{bmatrix}$  =  $(I + \Delta t A)^n (c_1 u_1 + c_2 u_2) = c_1 (1 + \Delta t \lambda_1)^n u_1 + c_2 (1 + \Delta t \lambda_2)$ 

![](_page_55_Picture_3.jpeg)

## *Observations: stiff systems*

- Whether the simulation blows up depends on $\left| \left( 1 + \Delta t \lambda_i \right)^n \right| \leq 1$ *n*  $\int (1 + \Delta t \lambda_i$ • If eigen values are of vast distinct sizes, a stiff system results
- The most negative eigen value determines the speed of simulation
- $\cdot$  It is also the part of the simulation that quickly dies out and not of interest

![](_page_56_Picture_4.jpeg)

*Try this!*

- 
- *❖* First order
	- $\Box$  convergent (e.g., a=d=- $5, c= d=0$
	- $a=d=0, b=2, d=-2)$
	- **<u>Q</u>** spiral combination of the above

 $\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$ 

**Second order** 

3, c=d=0)  
\n
$$
a=dd=0, b=2, d=-2
$$
\n
$$
a=1, c=1, d=1, d=0, b=2, d=0
$$
\n
$$
a=1, d=0, b=2, d=-2
$$
\n
$$
a=1, a=0, b=2, d=-2
$$
\n
$$
a=1, a=0, b=2, d=-2
$$
\n
$$
a=1, a=0, b=2, d=-2
$$
\n
$$
a=
$$

![](_page_57_Picture_9.jpeg)

 What is the difference between 1st and 2nd order simulations?

1st order 2nd order specify location directly

*y*

*x*

 $\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} =$ 

 $\ddagger$ 

specify velocity instead of location

![](_page_58_Figure_3.jpeg)

![](_page_58_Picture_4.jpeg)

- A simulation with  $a = d = -5$ ,  $b = c = 0$ 
	- $\Box$  can be free of oscillation if step size is properly chosen
	- **E** cannot be made oscillation free without a damping term

![](_page_59_Picture_3.jpeg)